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Influence of the water content in rock on the thermal neutron diffusion and diffusion cooling coefficients (by Monte Carlo simulations). II: – Quartz

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Abstract

The dependence of the thermal neutron diffusion parameters on the water content w in quartz, SiO₂, has been studied by means of Monte Carlo simulations of the pulsed neutron experiments for a number of series of samples. The density-removed diffusion cooling coefficient C^{M} varies hyperbolically between 39 400 000 and 4940 cm⁴s⁻¹(g/cm³)³ at the water content in the full range $0 \le w \le 1$. The obtained function $C^{\text{M}}(w)$ is compared with the analogous dependence for moisturized dolomite.

1. Introduction

Macroscopic parameters of the thermal neutron diffusion in a material depend on the scattering properties of the contributing substances. The diffusion coefficient, D, or the diffusion constant, D_0 , is roughly inversely proportional to the macroscopic total scattering cross-section, Σ_s . The diffusion cooling coefficient, C (with the correction, F), depends on Σ_s as well as on the scattering kernel (*i.e.* the differential scattering cross-section, $\Sigma_s(E' \rightarrow E)$, where E' and E are the thermal neutron energies before and after the collision). The relevant definitions, in which also the influence of the thermal neutron energy distribution is taken into account, can be found *e.g.* in [1] and [2]. A calculation of the diffusion cooling coefficient is always difficult (cf. [3], [4], [5]) and results for a wet rock material become uncertain. In this case, not only the values of the scattering cross-section, $\Sigma_s(E' \rightarrow E)$, in the two media is entirely dissimilar. Therefore, even a small change of the water content in a rock material results in a very significant change of the value of the diffusion cooling coefficient of the composed material.

An experimental procedure to determine the thermal neutron diffusion parameters for such complex media is usually most adequate. A pulsed neutron experiment (the so-called variable geometric buckling experiment) is then used. It is based on measurements of the time decay of the thermal neutron flux (after irradiation with the neutron burst) in a series of samples of the same material of varying size. The method is known in the experimental neutron physics, *e.g.* [6]. In the case of a rock material, some problem is created to keep a constant bulk density and homogeneity of samples. An additional serious technical problem appears when the experimental rock material has to have the water content precisely defined and repeatable in consecutive samples. The present-day computer techniques offer a possibility to simulate such experiments using Monte Carlo simulations of the neutron transport in the matter.

An influence of the water content on the thermal neutron diffusion parameters for one of the basic rock constituents, dolomite – $CaMg(CO_3)_2$, was investigated [7] with the Monte Carlo simulations of the variable buckling experiment. As said there, the dependences, which were found, cannot be treated as universal because the scattering properties of different elements are individual and usually very different. Here, we perform a study for another basic rock mineral, quartz – SiO₂. The idea of the simulated series of the experiments was presented

in detail in [7], [8], [9]. We remind here only the basic relationship necessary for interpretation of the experiments:

$$\lambda = \langle v\Sigma_{a} \rangle + D_{0} B^{2} - C B^{4} + F B^{6} - \dots , \qquad (1)$$

where λ is the decay constant of the fundamental mode of the time distribution of the thermal neutron pulsed flux, v is the thermal neutron speed, Σ_a is the macroscopic absorption crosssection, D_0 is the diffusion constant, C is the diffusion cooling coefficient, F is a correction term, and B^2 is the geometric buckling. It is defined [6] by the shape and size of the sample (including the extrapolation length) and for basic geometries is expressed by simple formulae [10]. Here, the spherical geometry of the experiment is kept and the buckling B^2 is defined as in [9] with a comment on the extrapolation length given in [7].

A fit of the function $\lambda(B^2)$ to the 'experimental' data λ_i obtained for different samples (different B_i^2) at the given water content, w, determines the values of the thermal neutron diffusion parameters, D_0 , C, F. The absorption rate, $\langle v \Sigma_a \rangle$, can be calculated with a high precision (cf. [1], [6]) and set into Eq.(1) as the known constant. A repetition of the simulations at the varying water content allows us to find the respective functions, $D_0 = D_0(w)$, C = C(w), and F = F(w).

2. Thermal neutron diffusion parameters of the contributing pure media

In the thermal neutron energy region (about 10^{-3} to 1 eV) the microscopic scattering cross-sections σ_s of quartz and of water differ not only in their values but also in the type of the energy dependence:

quartz, SiO₂
$$\sigma_s(E) \approx const. = \sigma_{sf}, \quad \mu(E) = const,$$
 (2a)
water, H₂O $\sigma_s(E) = f_s(E), \quad \mu(E) = f_{\mu}(E),$ (2b)

where μ is the average cosine of the scattering angle, and σ_{sf} is the microscopic scattering cross-section of the molecule built of atoms treated as free. The relevant macroscopic cross-sections are then defined by

$$\Sigma_{sf} = \Sigma_{s}(v_{0}) = 0.2541 \,\mathrm{cm}^{-1}, \qquad \text{for quartz } (\rho = 2.65 \,\mathrm{g/cm}^{3}) \qquad (3a)$$

and

$$\Sigma_{\rm s\,f} = 1.4921 \,\,{\rm cm^{-1}},$$

$$\Sigma_{\rm s}(v_0) \approx 3.98 \,\,{\rm cm^{-1}},$$

$$\langle \mu(E) \rangle \approx 0.2 \,\, (3b)$$

The data in Eqs (3) are based on the microscopic cross-sections from [11], and $v_0 = 2200$ m/s is the most probable thermal neutron velocity, corresponding to the energy $E \approx 0.0253$ eV.

Due to the characteristics summarized in Eqs (2) and (3), the thermal neutron energy--averaged diffusion parameters, D_0 , C, F, of the two media are strongly different. They were determined [7], [9] with the same Monte Carlo simulation method as mentioned above and are here quoted in Tables 1a and 2a. Cases marked as (i) contain results of a more accurate (on the order of B^6) fit of Eq.(1). Results given in Cases (ii) correspond to a fit of Eq.(1) with the accuracy of $O(B^4)$. An influence of the neglected correction F is then visible as a change of values of the parameters D_0 and C. They are sometimes helpful in such theoretical consideration of the neutron transport in which the correction F is not present explicitly.

Tables 1b and 2b contain the so-called density-removed equivalents of the thermal neutron pulsed parameters:

$$\langle v\Sigma_{a}\rangle^{M} = \langle v\Sigma_{a}\rangle \rho^{-1}, \qquad D_{0}^{M} = \rho D_{0}, \qquad C^{M} = \rho^{3} C, \qquad F^{M} = \rho^{5} F.$$
 (4)

Formulae (4) were obtained on the basis of theoretical definitions of the thermal neutron absorption-diffusion macroscopic parameters [1], [2], [6], [12]. More information on a reason for determination the density-removed parameters can be found in [7], [9], [12], [13].

SiO₂

Table 1a. Thermal neutron diffusion parameters of quartz ($\rho = 2.65 \text{ g/cm}^3$).

Case	$\langle v \Sigma_{a} \rangle$ [s ⁻¹]	D_0 [cm ² s ⁻¹]	C $[cm^4s^{-1}]$	$\frac{F}{[\mathrm{cm}^{6}\mathrm{s}^{-1}]}$
(i)		308 500	2 117 000	7 060 000
	1001	$\pm 2 \ 400$	$\pm 73~000$	±530000
(ii)	± 18	271 000	1 094 000	
		± 3700	$\pm 44~000$	

Case	$\langle \nu \Sigma_a \rangle^{\mathrm{M}}$ [s ⁻¹ /(g/cm ³)]	D_0^{M} [cm ² s ⁻¹ (g/cm ³)]	$\frac{C^{M}}{[\ cm^{4}s^{-1}(g/cm^{3})^{3}]}$	F^{M} [cm ⁶ s ⁻¹ (g/cm ³) ⁵]
(i)		817 500	39 400 000	922 000 000
	378	$\pm 6~400$	$\pm 1\ 360\ 000$	$\pm\ 68\ 000\ 000$
(ii)	± 7	718 100	20 350 000	
		± 9800	±830000	

Table 1b. Density-removed thermal neutron parameters of quartz.

H_2O

Table 2a. Thermal neutron diffusion parameters of water at 20 °C ($\rho = 0.99762$ g/cm³).

Case	$\langle v \Sigma_{a} \rangle$ [s ⁻¹]	$D_0 \ [cm^2 s^{-1}]$	C $[cm^4s^{-1}]$	F[cm ⁶ s ⁻¹]
(i)		35 450	4970	1530
	4882	± 150	± 580	± 520
(ii)	± 10	35 064	3340	
		± 89	±130	

 Table 2b.
 Density-removed thermal neutron parameters of water.

Case	$\langle \nu \Sigma_a \rangle^M$ [s ⁻¹ /(g/cm ³)]	D_0^{M} [cm ² s ⁻¹ (g/cm ³)]	$\frac{C^{M}}{[\ cm^{4}s^{-1}(g/cm^{3})^{3}]}$	F^{M} [cm ⁶ s ⁻¹ (g/cm ³) ⁵]
(i)		35 360	4940	1510
	4894	± 150	± 570	± 510
(ii)	± 10	34 981	3320	
		± 89	± 130	

3. Determination of the thermal neutron diffusion parameters for quartz containing water

The variable buckling experiment has been simulated at various water contents in quartz samples, in the same way as described in [7]. The thermal neutron transport parameters, D_0 , C, F, have been determined in each case. The thermal neutron absorption rate, $\langle v\Sigma_a \rangle$, can be calculated exactly from the elemental composition and the microscopic absorption cross--sections of the contributing elements, and has been always used in the fitting procedure as the known constant. The results of the individual simulated series of experiments (*i.e.* the decay constants λ_i and the neutron parameters obtained finally) are collected below. The same nomenclature as in [7] is kept here:

Moisturized quartz – material which contains quartz, SiO_2 , and an amount of water, *Water content* – mass contribution of water in the moisturized material, *i.e.* the ratio

of

the mass of water to the mass of quartz+water, given as

w – weight fraction, or as

p – weight per cent.

The experiments have been simulated for the water contents equal to 1, 2, 4, 6, 8, 10, and 20 per cent.

3.1. Moisturized quartz, p = 1 %

The material density is $\rho = 2.633 \text{ g/cm}^3$. The thermal neutron decay constants λ_i obtained from the simulations for individual spheres are listed in Table 3 (where R_g is the geometric radius of the sphere, and $\sigma(\lambda)$ is the standard deviation of the λ value determined). The resulting plot $\lambda = \lambda(B^2)$ is shown in Fig. 1. The corresponding thermal neutron diffusion parameters from the fits (i) and (ii), defined in paragraph 2, are collected in Table 4a. The related density-removed parameters are given in Table 4b.

The same scheme of presentation of the results is used for all following simulations of the pulsed thermal neutron flux at the varying water content in quartz.

R _g	λ	$R_{\rm g}$	λ	$R_{\rm g}$	λ
C C	σ(λ)	6	σ(λ)		$\sigma(\lambda)$
[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$
5.2	27 267	6.5	22 633	10.0	14 336
	150		88		52
5.5	26 153	7.0	21 195	13.0	10 406
	123		199		45
6.0	24 134	8.0	18 388	20.0	5 865
	104		31		18

Table 3. Decay constants λ obtained from the simulated experiment for quartz with the 1 % water content.

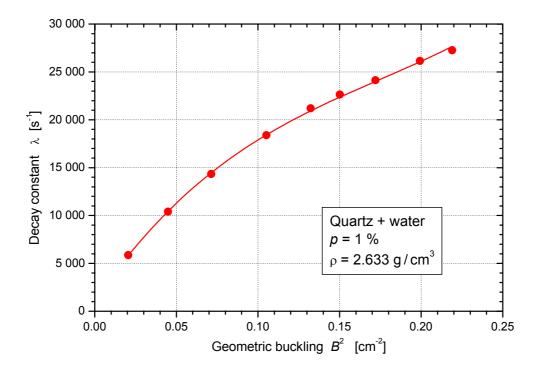


Fig. 1. Results of the simulations of the variable buckling experiment for quartz with the 1 % water content.

Table 4a. Thermal neutron diffusion parameters determined for quartz with the 1 % water content ($\rho = 2.633 \text{ g/cm}^3$).

			Fixed		Fitted	
Case	Range of $R_{\rm g}$	Range of B^2	$\langle v \Sigma_{a} \rangle$	D_0	C	F
	[cm]	$[cm^{-2}]$	$[s^{-1}]$	$[cm^2s^{-1}]$	$[cm^4s^{-1}]$	$[cm^{6}s^{-1}]$
(i)		0.021 . 0.210		249 000	1 004 000	1 920 000
	$5.2 \div 20$	$0.021 \div 0.219$	1114	$\pm 2\ 400$	$\pm 39\ 000$	$\pm 150\ 000$
(ii)	$5.2 \div 20$	0.022.0220	±17	214 700	477 000	
		$0.022 \div 0.229$		$\pm 6\ 600$	$\pm 46\ 000$	

Table 4b. Density-removed thermal neutron parameters of quartz with the 1 % water content.

Case	$\langle v \Sigma_a \rangle^M$ [s ⁻¹ /(g/cm ³)]	D_0^{M} [cm ² s ⁻¹ (g/cm ³)]	C^{M} [cm ⁴ s ⁻¹ (g/cm ³) ³]	$F^{\rm M}$ [cm ⁶ s ⁻¹ (g/cm ³) ⁵]
(i)		655 500	18 320 000	243 000 000
	423	$\pm 6~400$	\pm 720 000	$\pm 19\ 000\ 000$
(ii)	± 6	565 000	8 700 000	
		$\pm 17\ 000$	±840000	

3.2. Moisturized quartz, p = 2 %

The material density is $\rho = 2.617 \text{ g/cm}^3$.

122

85

27 470

5.5

water content.					
R _g	λ	Rg	λ	R _g	λ
C	σ(λ)	C	σ(λ)	C C	σ(λ)
[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$	[cm]	[s ⁻¹]
4.5	32 635	6.0	24 990	10.0	13 833
	153		103		21
4.7	31 454	6.5	23 014	13.0	9 792
	104		86		18
5.0	29 805	7.0	21 164	20.0	5 436

8.0

36

35

18 183

10

Table 5. Decay constants λ obtained from the simulated experiment for quatrz with the 2 % water content.

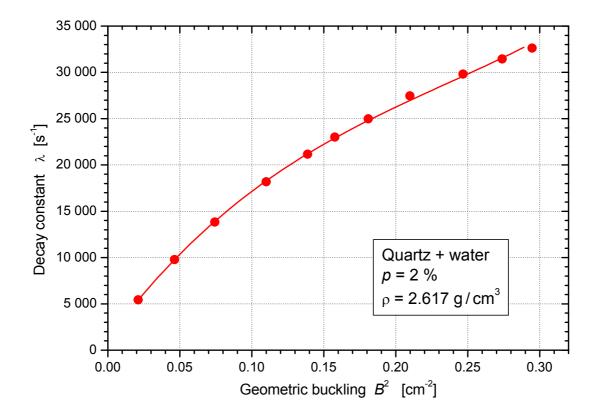


Fig. 2. Results of the simulations of the variable buckling experiment for quartz with the 2 % water content.

			Fixed		Fitted	
Case	Range of <i>R</i> g [cm]	Range of B^2 [cm ⁻²]	$\langle v \Sigma_{a} \rangle$ [s ⁻¹]	$D_0 \ [m cm^2 s^{-1}]$	C $[cm^4s^{-1}]$	F[cm ⁶ s ⁻¹]
(i)		0.021 ÷ 0.295		210 100	593 000	840 000
	$4.5 \div 20$	$0.021 \div 0.293$	1225	± 1500	$\pm 21\ 000$	$\pm 64\ 000$
(ii)	$4.3 \div 20$	0.021 . 0.204	±17	190 200	313 000	
		$0.021 \div 0.304$		± 3800	$\pm 23\ 000$	

Table 6a. Thermal neutron diffusion parameters determined for quartz with the 2 % water content ($\rho = 2.617 \text{ g/cm}^3$).

Table 6b. Density-removed thermal neutron parameters of quartz with the 2 % water content.

Case	$\langle v \Sigma_a \rangle^M$ [s ⁻¹ /(g/cm ³)]	D_0^{M} [cm ² s ⁻¹ (g/cm ³)]	C^{M} [cm ⁴ s ⁻¹ (g/cm ³) ³]	F^{M} [cm ⁶ s ⁻¹ (g/cm ³) ⁵]
(i)		549 700	10 620 000	103 100 000
	468	± 3900	±380000	$\pm\ 7\ 900\ 000$
(ii)	± 6	498 000	5 600 000	
		$\pm 10\ 000$	$\pm400\;000$	

3.3. Moisturized quartz, p = 4 %

The material density is $\rho = 2.584 \text{ g/cm}^3$.

Table 7. Decay constants λ obtained from the simulated experiment for quartz with the 4 % water content.

Rg	λ	R _g	λ	R _g	λ
C	σ(λ)		σ(λ)	C	$\sigma(\lambda)$
[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$
4.5	33 888	6.5	22 231	13.0	8 715
	96		69		17
5.0	30 189	7.0	20 206	20.0	4 874
	164		23		13
5.5	27 089	8.0	16 996		·
	62		21		
6.0	24 421	10.0	12 570		
	51		27		

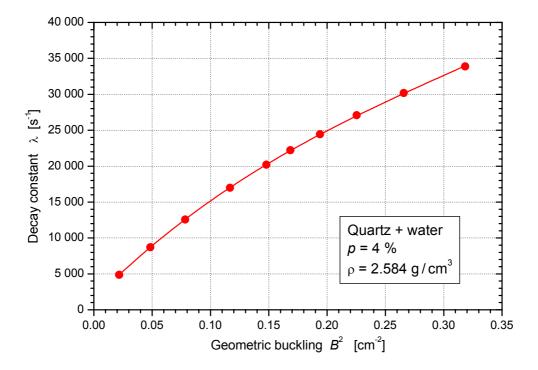


Fig. 3. Results of the simulations of the variable buckling experiment for quartz with the 4 % water content.

Table 8a. Thermal neutron diffusion parameters determined for quartz with the 4 % water content ($\rho = 2.584 \text{ g/cm}^3$).

			Fixed		Fitted	
Case	Range of <i>R</i> g [cm]	Range of B^2 [cm ⁻²]	$\langle v \Sigma_{a} \rangle$ [s ⁻¹]	$D_0 \ [m cm^2 s^{-1}]$	C $[cm^4s^{-1}]$	F [cm ⁶ s ⁻¹]
(i)		$0.022 \div 0.318$		163 210	292 400	316 000
	$4.5 \div 20$	$0.022 \div 0.518$	1443	± 450	$\pm 5\ 200$	\pm 14 000
(ii)	$4.3 \div 20$	$0.022 \div 0.325$	±16	152 700	175 000	
		$0.022 \div 0.323$		$\pm 2\ 000$	$\pm \ 11 \ 000$	

Table 8b. Density-removed thermal neutron parameters of quartz with the 4 % water content.

Case	$\langle v \Sigma_a \rangle^M$ [s ⁻¹ /(g/cm ³)]	D_0^{M} [cm ² s ⁻¹ (g/cm ³)]	C^{M} [cm ⁴ s ⁻¹ (g/cm ³) ³]	F^{M} [cm ⁶ s ⁻¹ (g/cm ³) ⁵]
(i)		421 700	5 045 000	36 400 000
	558	$\pm 1\ 200$	$\pm 90\ 000$	$\pm 1\ 600\ 000$
(ii)	± 6	394 500	3 030 000	
		± 5 100	±190000	

3.4. Moisturized quartz, p = 6 %

The material density is $\rho = 2.551 \text{ g/cm}^3$.

Rg	λ	$R_{ m g}$	λ	Rg	λ
	σ(λ)	C C	$\sigma(\lambda)$	C C	σ(λ)
[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$
4.5	33 269	6.5	20 876	13.0	7 949
	134		45		12
5.0	29 206	7.0	18 870	20.0	4 544
	71		27		12
5.5	25 938	8.0	15 691		
	56		19		
6.0	23 211	10.0	11 484		
	48		16		

Table 9. Decay constants λ obtained from the simulated experiment for quartz with the 6 % water content.

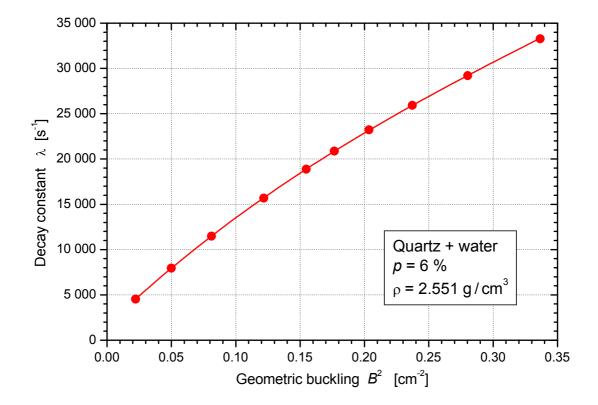


Fig. 4. Results of the simulations of the variable buckling experiment for quartz with the 6 % water content.

			Fixed		Fitted	
Case	Range of <i>R</i> g [cm]	Range of B^2 [cm ⁻²]	$\langle v \Sigma_{a} \rangle$ [s ⁻¹]	$D_0 \ [ext{cm}^2 ext{s}^{-1}]$	C $[cm^4s^{-1}]$	F[cm ⁶ s ⁻¹]
(i)		$0.022 \div 0.337$		133 890	169 500	152 500
	$4.5 \div 20$	$0.022 \div 0.557$	1655	± 240	$\pm 3\ 000$	$\pm 8 300$
(ii)	$4.3 \div 20$	$0.022 \div 0.340$	±16	129 550	116 400	
		$0.022 \div 0.340$		± 870	$\pm 5\ 000$	

Table 10a. Thermal neutron diffusion parameters determined for quartz with the 6 % water content ($\rho = 2.551 \text{ g/cm}^3$).

Table 10b. Density-removed thermal neutron parameters of quartz with the 6 % water content.

Case	$\langle v \Sigma_a \rangle^M$ [s ⁻¹ /(g/cm ³)]	D_0^{M} [cm ² s ⁻¹ (g/cm ³)]	C^{M} [cm ⁴ s ⁻¹ (g/cm ³) ³]	F^{M} [cm ⁶ s ⁻¹ (g/cm ³) ⁵]
(i)		341 550	2 814 000	16 470 000
	649	± 610	$\pm 50\ 000$	$\pm 900\ 000$
(ii)	± 6	330 500	1 933 000	
		$\pm 2 \ 200$	$\pm 83\ 000$	

3.5. Moisturized quartz, p = 8 %

The material density is $\rho = 2.518 \text{ g/cm}^3$.

Table 11. Decay constants λ obtained from the simulated experiment for quartz with the 8 % water content.

Rg	λ	R _g	λ	R _g	λ
_	σ(λ)	_	$\sigma(\lambda)$		$\sigma(\lambda)$
[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$
4.5	31 978	6.5	19 573	13.0	7 428
	35		19		26
5.0	27 900	7.0	17 654	20.0	4 377
	32		30		16
5.5	24 628	8.0	14 609		
	48		16		
6.0	21 870	10.0	10 641		
	38		13		

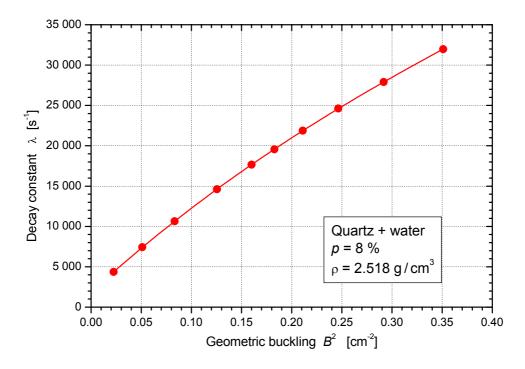


Fig. 5. Results of the simulations of the variable buckling experiment for quartz with the 8 % water content.

Table 12a. Thermal neutron diffusion parameters determined for quartz with the 8 % water content ($\rho = 2.518 \text{ g/cm}^3$).

			Fixed		Fitted	
Case	Range of <i>R</i> g [cm]	Range of B^2 [cm ⁻²]	$\langle v \Sigma_{a} \rangle$ [s ⁻¹]	D_0 [cm ² s ⁻¹]	C $[cm^4s^{-1}]$	F [cm ⁶ s ⁻¹]
(i)		0.023 ÷ 0.351		113780	105 600	73 900
	4.5 ÷ 20	$0.023 \div 0.331$	1861	± 170	± 1700	± 3700
(ii)	$4.3 \div 20$	$0.023 \div 0.354$	±15	110 120	72 300	
		$0.023 \div 0.334$		± 540	$\pm 2 200$	

Table 12b. Density-removed thermal neutron parameters of quartz with the 8 % water content.

Case	$\langle v \Sigma_a \rangle^M$ [s ⁻¹ /(g/cm ³)]	D_0^{M} [cm ² s ⁻¹ (g/cm ³)]	$\frac{C^{M}}{[cm^{4}s^{-1}(g/cm^{3})^{3}]}$	F^{M} [cm ⁶ s ⁻¹ (g/cm ³) ⁵]
(i)		286 490	1 686 000	7 480 000
	739	± 430	±27000	$\pm370~000$
(ii)	± 6	277 300	1 154000	
		$\pm 1 300$	$\pm 35\ 000$	

3.6. Moisturized quartz, p = 10 %

The material density is $\rho = 2.485 \text{ g/cm}^{-3}$.

Rg	λ	Rg	λ	Rg	λ
- C	σ(λ)		$\sigma(\lambda)$		$\sigma(\lambda)$
[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$
4.5	30 633	6.5	18 444	13.0	7 039
	50		27		18
5.0	26 575	7.0	16 588	20.0	4 290
	45		21		16
5.5	23 330	8.0	13 719		
	45		22		
6.0	20 664	10.0	10 023		
	43		26		

Table 13. Decay constants λ obtained from the simulated experiment for quartz with the 10 % water content.

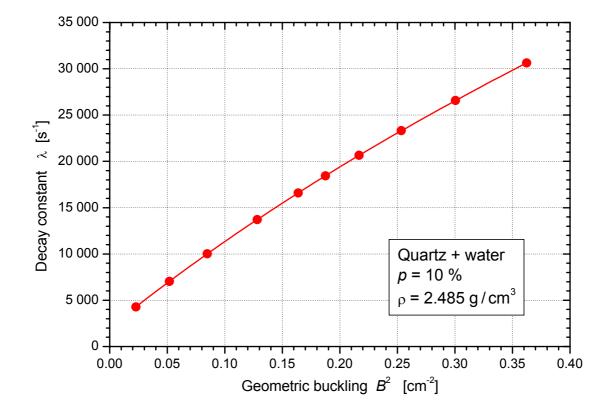


Fig. 6. Results of the simulations of the variable buckling experiment for quartz with the 10 % water content.

			Fixed		Fitted	
Case	Range of <i>R</i> g [cm]	Range of B^2 [cm ⁻²]	$\langle v \Sigma_{a} \rangle$ [s ⁻¹]	$D_0 \ [\mathrm{cm}^2 \mathrm{s}^{-1}]$	C $[cm^4s^{-1}]$	F[cm ⁶ s ⁻¹]
(i)		$0.023 \div 0.362$		99 730	74 600	47 000
	$4.5 \div 20$	$0.023 \div 0.302$	2061	± 150	$\pm 1 400$	± 2900
(ii)	$4.3 \div 20$	$0.023 \div 0.364$	±15	97 280	53 000	
		$0.023 \div 0.304$		± 390	± 1600	

Table 14a. Thermal neutron diffusion parameters determined for quartz with the 10 % water content ($\rho = 2.485 \text{ g/cm}^3$).

Table 14b. Density-removed thermal neutron parameters of quartz with the 10 % water content.

Case	$\langle v \Sigma_a \rangle^M$ [s ⁻¹ /(g/cm ³)]	D_0^{M} [cm ² s ⁻¹ (g/cm ³)]	C^{M} [cm ⁴ s ⁻¹ (g/cm ³) ³]	F^{M} [cm ⁶ s ⁻¹ (g/cm ³) ⁵]
(i)		247 820	1 144 000	4 450 000
	829	± 360	$\pm 21\ 000$	$\pm270\;000$
(ii)	± 6	241 740	814 000	
		± 960	$\pm 25\ 000$	

3.7. Moisturized quartz, p = 20 %

The material density is $\rho = 2.320 \text{ g/cm}^3$.

Table 15. Decay constants λ obtained from the simulated experiment for quartz with the 20 % water content.

R _g	λ	R _g	λ	R _g	λ
	$\sigma(\lambda)$		$\sigma(\lambda)$	C C	$\sigma(\lambda)$
[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$	[cm]	$[s^{-1}]$
4.5	25 353	6.5	14 929	13.0	6 336
	60		14		18
5.0	21 764	7.0	13 480	20.0	4 4 3 0
	33		24		12
5.5	18 942	8.0	11 238		
	23		20		
6.0	16 728	10.0	8 471		
	22		25		

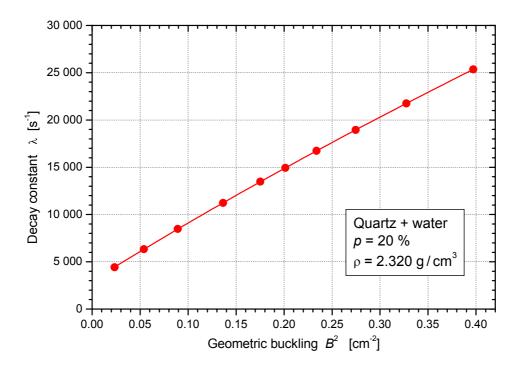


Fig. 7. Results of the simulations of the variable buckling experiment for quartz with the 20 % water content.

Table 16a. Thermal neutron diffusion parameters determined for quartz with the 20 % water content ($\rho = 2.320 \text{ g/cm}^3$).

			Fixed		Fitted	
Case	Range of <i>R</i> g [cm]	Range of B^2 [cm ⁻²]	$\langle v \Sigma_{a} \rangle$ [s ⁻¹]	$D_0 \ [m cm^2 s^{-1}]$	C $[cm^4s^{-1}]$	F[cm ⁶ s ⁻¹]
(i)		$0.023 \div 0.397$		63 360	21 600	10 000
	4.5 ÷ 20	$0.023 \div 0.397$	2972	± 170	$\pm 1 400$	± 2700
(ii)	$4.3 \div 20$	0.022 . 0.208	±13	62 770	16 820	
		$0.023 \div 0.398$		± 120	± 480	

Table 16b. Density-removed thermal neutron parameters of quartz with the 20 % water content.

Case	$\langle v \Sigma_a \rangle^M$ [s ⁻¹ /(g/cm ³)]	D_0^{M} [cm ² s ⁻¹ (g/cm ³)]	C^{M} [cm ⁴ s ⁻¹ (g/cm ³) ³]	F^{M} [cm ⁶ s ⁻¹ (g/cm ³) ⁵]
(i)		146 990	270 000	670 000
	1281	± 390	$\pm 17\ 000$	$\pm 180\ 000$
(ii)	± 6	145 620	210 000	
		± 280	$\pm 6\ 000$	

4. Dependence of the thermal neutron diffusion coefficients on the water content in quartz

4.1. The diffusion constant D_0 and the diffusion coefficient D

A dependence of the thermal neutron diffusion constant on the water content in quartz, $D_0(w)$, obtained from the simulations performed, is shown in Table 17 and in Fig. 8. The corresponding values of the diffusion coefficient, D, can be found from the relation:

$$\langle D(E)\rangle(w) = \left\langle \frac{1}{v} \right\rangle D_0(w) = \frac{\sqrt{\pi}}{2v_0} D_0(w) .$$
(5)

The diffusion constants D_{0C} have been independently calculated from an approximate theoretical formula for homogeneous mixtures of the hydrogenous and non-hydrogenous components [2], [14]. The data in Table 17 and the plot in Fig. 8 show that values of the diffusion constant obtained from the simulations, D_0 , and from the approximate calculation, D_{0C} , are very close.

			1
W	ρ	$\begin{bmatrix} D_0 \\ cm^2 s^{-1} \end{bmatrix}$	D_{0C} [cm ² s ⁻¹]
	[g/cm ³]		
0.00	2.65	308 500	~ 333 030
		$\pm 2\ 400$	± 430
0.01	2.633	249 000	$\sim 267\ 100$
		$\pm 2 400$	$\pm 2 300$
0.02	2.617	210 100	~ 223 400
		$\pm 1 500$	$\pm 3\ 200$
0.04	2.584	163 210	~ 169 200
		± 450	$\pm 3 600$
0.06	2.551	133 890	~ 136 800
		± 240	± 3500
0.08	2.518	113 780	~ 115 400
		± 170	$\pm 3\ 300$
0.10	2.485	99 730	~ 100 100
		± 150	$\pm 3\ 100$
0.20	2.320	63 360	$\sim 62\ 400$
		± 170	$\pm 2\ 200$
1.00	0.99762	35 450	~ 33 400
		± 150	$\pm 1 400$

Table 17. Dependence of the fitted, D_0 , and calculated, D_{0C} , diffusion constants of moisturized quartz on the water content.

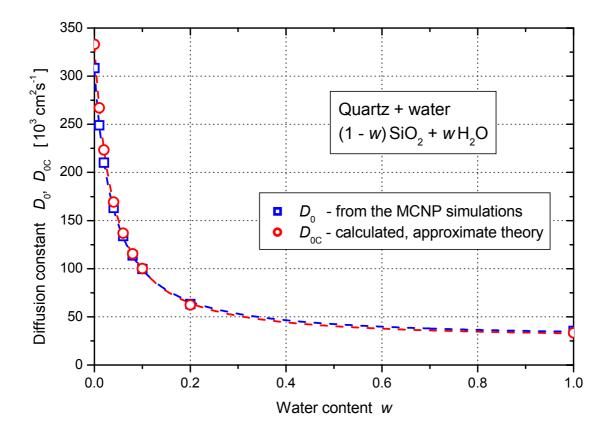


Fig. 8. Thermal neutron diffusion constant of moisturized quartz as a function of the water content.

4.1. The diffusion cooling coefficient C with its correction F

The dependence of the density-removed thermal neutron diffusion parameters, D_0^M , C^M , F^M , on water content *w*, obtained from the performed Monte Carlo simulations, is presented in Table 18.

The density-removed diffusion cooling coefficient C^{M} as a function of the water content in moisturized quartz is plotted in Fig. 9.

		М	М	М
W	ρ	D_0^{M}	C^{M}	F^{M}
	$[g/cm^3]$	$[cm^2 s^{-1} (g/cm^3)]$	$[cm^4 s^{-1} (g/cm^3)^3]$	$[cm^6 s^{-1} (g/cm^3)^5]$
0.00	2.65	817 500	39 400 000	922 000 000
		$\pm 6\ 400$	$\pm 1 \ 360 \ 000$	$\pm \ 68\ 000\ 000$
0.01	2.633	655 500	18 320 000	243 000 000
		$\pm 6\ 400$	\pm 720 000	$\pm 19\ 000\ 000$
0.02	2.617	549 700	10 620 000	103 100 000
		± 3900	\pm 380 000	$\pm 7 \ 900 \ 000$
0.04	2.584	421 700	5 045 000	36 400 000
		$\pm 1\ 200$	$\pm 90\ 000$	$\pm 1\ 600\ 000$
0.06	2.551	341 550	2 814 000	16 470 000
		± 610	$\pm 50\ 000$	$\pm 900\ 000$
0.08	2.518	286 490	1 686 000	7 480 000
		± 430	$\pm27\;000$	$\pm370~000$
0.10	2.485	247 820	1 144 000	4 450 000
		± 360	$\pm 21\ 000$	$\pm270\;000$
0.20	2.320	146 990	270 000	670 000
		± 390	$\pm 17\ 000$	$\pm\ 180\ 000$
1.00	0.99762	35 360	4 940	1 510
		±150	± 570	± 510

Table 18. Dependence of the density-removed thermal neutron diffusion parameters of moisturized quartz on the water content.

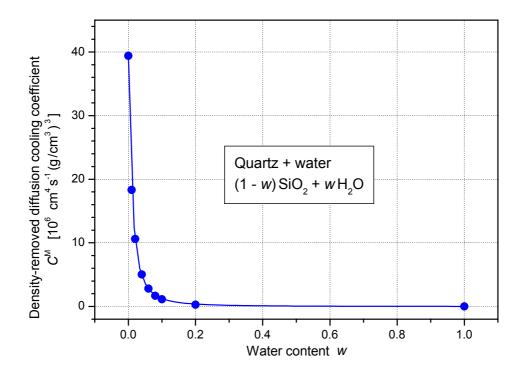


Fig. 9. Density-removed diffusion cooling coefficient C^{M} of moisturized quartz as a function of the water content.

The dependence $C^{M}(w)$ has a hyperbolic shape, as found in [7]. The function

$$C^{\rm M}(w) = a (w + w_0)^{b}$$
(6)

has been fitted to the 'experimental' data C_i^M and the following parameters have been obtained:

$$\begin{array}{l} w_0 = 0.02125 \pm 0.00090 \\ a = 19500 \pm 2900 \text{ cm}^4 \text{s}^{-1} (\text{g/cm}^3)^3 \\ b = -1.977 \pm 0.060 \end{array} \right\} .$$
 (6a)

Formula (6) is slightly different than that reported in [7] for dolomite. When additional simulations for moisturized dolomite were made for very low water contents, the function (6) was found to be better and, therefore, it has been also used for moisturized quartz here.

5. Conclusions

The formulae $C^{M}(w)$ obtained for moisturized quartz and dolomite [7] are accurate for $w \in [0, 0.1]$, and can be used approximately up to w = 0.2. At higher water contents (which are not interested for geophysical interpretation of neutron log measurements) the formulae would be inaccurate as the simulations for $w \in (0.2, 1.0)$ have not been done. The revised parameters for dolomite, according to the fit (6), are

$$w_{0} = 0.0356 \pm 0.0050$$

$$a = 11900 \pm 5500 \text{ cm}^{4} \text{s}^{-1} (\text{g/cm}^{3})^{3}$$

$$b = -2.20 \pm 0.23$$
(6b)

The dependences $C^{M}(w)$ obtained for quartz and dolomite are, as expected, similar but the values at a particular water content are significantly different. Moreover, the quartz-to-dolomite diffusion cooling ratio, $C_{Q}^{M}(w)/C_{D}^{M}(w)$, is not constant as a function of the water content. The ratio is plotted in Fig. 10.

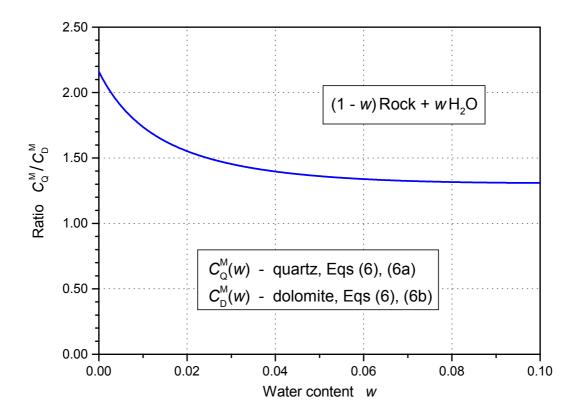


Fig. 10. Ratio of the diffusion cooling coefficients C^{M} of moisturized quartz (Q) and dolomite (D) as a function of the water content w.

The obtained results show that the influence of the water content in the rock on the thermal neutron diffusion cooling properties is individual for different rocks. Therefore, similar simulations as performed here for quartz and in [7] for dolomite are also desired at least for the third basic rock mineral, calcite – $CaCO_3$.

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